

time span simulated (fig. S6 and table S8). These results, combined with the lack of resolution within superclades of the metazoan tree, argue against models of metazoan radiation in which the temporal window of diversification is much larger (48).

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the symmetry representations correspond to even or odd parity of an energy level.

We now report the successful use of LID to deplete the population of the  $B_{2u}$  isomer in a sample of gaseous ethylene, followed by monitoring of the subsequent spin conversions for the return to equilibrium. We measured isomer concentrations by recording the absorption intensities of spectral lines with appropriate  $J$ ,  $K_a$ , and  $K_c$  quantum numbers. Our experimental setup uses two  $\text{CO}_2$  lasers (Edinburgh Instruments PL3 as the separation laser and a home-built laser as the probe) and three glass cells (for separation, test, and reference) (16). We measured the spin conversion rates for  $^{13}\text{CH}_3\text{F}$  with this setup and obtained good agreement with the published results (6, 7).

For the ethylene study, the experimental schemes are shown in Table 2, where the reported results from high-resolution infrared spectroscopy (17) were used to calculate the frequency offsets between the  $\text{C}_2\text{H}_4$  transition frequencies and the  $\text{CO}_2$  laser frequencies. Application of the LID technique for the separation of nuclear spin isomers requires that a molecular transition be near-coincident with a  $\text{CO}_2$  laser line. Here, the 10P44 laser line with a power of 6 W was used. Its frequency was tuned about 20 MHz above the center frequency by adjusting the laser cavity length to set it in the red wing of the  $9_{0,9} \leftarrow 10_{1,9}$  line of the  $\nu_7$  band of ethylene. This frequency selectively excited the  $B_{2u}$  isomer, with the other three isomers using as a buffer gas. The  $B_{2u}$  molecules drift, by the LID effect, along the direction of the separation laser beam in the separation cell, thereby depleting the  $B_{2u}$  species and enriching the  $A_g$ ,  $B_{1g}$ , and  $B_{3u}$  species at the entrance end of the cell; this direction of drift corresponds to an increase in the collision cross section upon excitation. The nonequilibrium population was then transferred through a valve from the near end of the separation cell to the test cell. For high sensitivity, we measured differential absorption by splitting the probe beam to acquire simultaneous data from the test cell and the reference cell with a population at thermal equilibrium. We determined normalized absorption intensity differences for appropriate probe lines to

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populations. Very similar signals were also observed for alternative  $B_{2u}$  and  $B_{3u}$  probe resonances (cases 2 and 3 in Table 2). We tried to monitor the  $B_{1g}$  population dynamics but were not successful because the line intensity of the resonant  $26_{10,16} \leftarrow 27_{9,18}$  transition was too weak. The signals in the third period show the relaxation due to the conversion among spin isomers. A model function  $A \exp(-\gamma t) + B$  (where  $A$  is the integrated intensity,  $\gamma$  is the observed conversion rate constant, and  $B$  is the baseline offset) was fitted to the decay data of Fig. 1 to give the solid smooth curve shown with a rate constant  $\gamma = 8.09 (\pm 0.10) \times 10^{-4} \text{ s}^{-1}$ .

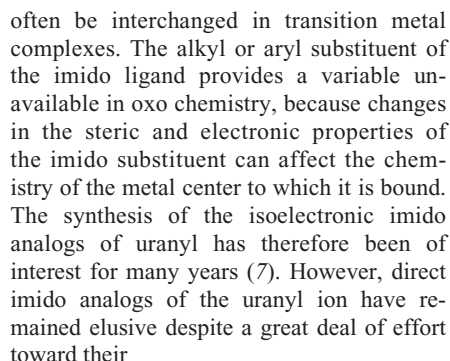
The data clearly show that the concentration of the  $A_g$  species is almost constant in time, whereas monoexponential kinetics are observed for recovery of the depleted  $B_{2u}$  population and decay of the enriched  $B_{3u}$  population. Furthermore, the  $B_{2u}$  signal does not return to the original zero-difference baseline, and the  $B_{3u}$  signal overshoots the baseline and asymptotically approaches a new equilibrium level. These general phenomena can be qualitatively explained using Curl's theory of state mixing (19). We assume that conversion of nuclear spin isomers of  $\text{C}_2\text{H}_4$  is allowed between the  $B_{2u}$  and  $B_{3u}$  isomers, and between the  $A_g$  and  $B_{1g}$  isomers, but forbidden between species of opposite inversion symmetry. Specifically, molecular "doorway" states are posited, between either  $B_{2u}$  and  $B_{3u}$  or  $A_g$  and  $B_{1g}$ , that are so close in energy that the weak intramolecular nuclear spin-rotation and spin-spin interactions of  $\text{C}_2\text{H}_4$  can induce mixing between them. This mixing is interrupted by collisions, which promote interconversion between either the  $B_{2u}$  and  $B_{3u}$  or the  $A_g$  and  $B_{1g}$

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Te40-14x18.3 (e40-14r)r

8110310311-1.1359.1.630.1(s3.17734 Td [(imid)Replac73.5 8 1 Tf90.4995)-8944-7



The imido ligand ( $\text{NR}^{2-}$ ) is isoelectronic with the oxo ligand, and the two groups can



**Separation and Conversion Dynamics of Four Nuclear Spin Isomers of Ethylene**

Zhen-Dong Sun, Kojiro Takagi and Fusakazu Matsushima  
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Editor's Summary

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